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SUSY R-parity violating contributions to the width differences for $D - \bar{D}$ and $B_{d,s} - \bar{B}_{d,s}$ systems

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Abstract: We study R-parity violating contributions to the mixing parameter $y$ for $D^0 - \bar{D}^0$ and $B^0_{d,s} - \bar{B}^0_{d,s}$ systems. We first obtain general expressions for new physics contributions to $y$ from effective four fermion operators. We then use them to study R-parity contributions. We find that R-parity violating contributions to $D^0 - \bar{D}^0$ mixing, and $B^0_{d,s} - \bar{B}^0_{d,s}$ to be small. There may be sizable contribution to $B^0_{d,s} - \bar{B}^0_{d,s}$ mixing. We also obtain some interesting bounds on R-parity violating parameters using known Standard Model predictions and experimental data.

Keywords: Supersymmetry Phenomenology, Beyond Standard Model.

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1. Introduction

Mixing between a neutral meson with specific flavor and its anti-meson provides powerful test for the Standard Model (SM) and new physics (NP) beyond. Mixing has been observed in $K^0 - \bar{K}^0$, $B^0_d - \bar{B}^0_d$, and $B^0_s - \bar{B}^0_s$. Evidence for $D^0$ mixing has been recently reported by the BaBar [2] and Belle [3, 4] collaborations.

Two parameters, $x = \Delta M/\Gamma$ and $y = \Delta \Gamma/2\Gamma$, are often used to describe the mixing between a meson and its anti-meson. Here $\Gamma$ is the life-time of the meson. $\Delta M = m_2 - m_1$, $\Delta \Gamma = \Gamma_2 - \Gamma_1$ with “1” and “2” indicating the CP odd and CP even states, respectively, in the limit of CP conservation. $\Delta M$ and $\Delta \Gamma$ are related to the mixing matrix elements $M_{12}$ and $\Gamma_{12}$ in the Hamiltonian by

$$\Delta M - i\Delta \Gamma/2 = 2\sqrt{(M_{12} - i\Gamma_{12}/2)(M_{12}^* - i\Gamma_{12}^*/2)}.$$ 

If a new particle has flavor changing neutral current (FCNC) interaction, a non-zero contribution to $M_{12}$ can be easily generated by exchanging this new particle in the intermediate state, tree or loop. The parameter $\Gamma_{12}$ must come from the absorptive part which requires the intermediate states be light degrees of freedom to whom the meson can decay into. This fact severely constrains the contributions to $\Gamma_{12}$ from NP. Due to this reason there is less theoretical work on new physics contributions to $\Gamma_{12}$ than that for $M_{12}$. In this work, we study the $\Gamma_{12}$ parameter in the present of NP, taking SUSY R-parity violating (RPV) interaction as an explicit example.
There are three types of R-Parity violating (RPV) terms \[3]\: 
\[
\frac{\lambda_{ijk}}{2} L_i^L L_i^L E_R^k , \quad \frac{\lambda_{ijk}^*}{2} L_i^L Q_i^L D_R^k , \quad \frac{\lambda_{ijk}^{"}}{2} U_R^\gamma D_R^k D_R^k , 
\]  
(1.1)
where \(i, j\) and \(k\) are the generation indices; \(L_L, Q_L, E_R, D_R\) and \(U_R\) are the chiral superfields which transform under the SM gauge group \(SU(3)_C \times SU(2)_L \times U(1)_Y\) as \(L_L: (1, 2, -1),\ E_R: (1, 1, -2),\ Q_L: (3, 2, 1/3),\ U_R: (3, 1, 4/3)\ and \(D_R: (3, 1, -2/3)\). The charge conjugated field \(\psi_R^*\) is defined as \(\psi_R^* = C(\bar{\psi}_R)^T\). We will consider each of these R-parity contributions to \(\Delta \Gamma_{12}\) for meson mixing separately. In that case, as the term proportional to \(\lambda_{ijk}\) involves only leptons, it will not contribute to meson mixing, since we are not considering pairs of \(\lambda_{ijk}\) and \(\lambda_{ijk}^*\) couplings to be non-zero at the same time. We only need to consider the last two terms up to one loop level.

At the tree level, we have the following terms relevant to us by exchange s-fermions,

\[
\mathcal{L}_{\text{eff}}(\lambda') = \frac{\lambda_{ijk}}{2m^2_{\tilde{d}_L}} \bar{d}_{L_i} \gamma^\mu u_{L_i} \bar{u}_{L_j} \gamma^\mu u_{L_j} - \frac{\lambda_{ijk}^*}{2m^2_{\tilde{e}_L}} \bar{e}_{L_i} \gamma^\mu u_{L_i} \bar{d}_{L_j} \gamma^\mu d_{L_j} - \frac{\lambda_{ijk}^{"}}{2m^2_{\tilde{d}_L}} \bar{d}_{L_i} \gamma^\mu u_{L_i} \bar{d}_{L_j} \gamma^\mu d_{L_j} \gamma^\mu d_{R_k} \gamma^\mu d_{R_k}^\prime \gamma^\mu d_{R_k}^\prime ,
\]

(1.2)

The first two terms in \(\mathcal{L}_{\text{eff}}(\lambda'^*)\) contribute to \(\Gamma_{12}\) for \(D^0 - \bar{D}^0\) mixing. Except the first term in \(\mathcal{L}_{\text{eff}}(\lambda')\), all terms contribute to \(B^0_{d,s} - \bar{B}^0_{d,s}\) mixing.

It is clear that from the above Lagrangian at the tree level, non-zero \(M_{12}\) can be generated. Constraints have been obtained using \(\Delta M\) for various meson mixing. However, in order to generate a non-zero \(\Gamma_{12}\) additional loop corrections are needed from the above four fermion interactions.

There are short and long distance contributions to \(y\) or \(\Gamma_{12}\). The calculations for long distance contributions are very difficult to handle due to our poor understanding of QCD at low energies. It is expected that long distance contributions become less and less important when energy scale becomes higher and higher, and perturbative short distance contributions will become the dominant one. We therefore will restrict ourselves to mesons containing a heavy c or b quark and to study the short distance contributions \(\Gamma_{12}\) for \(D^0 - \bar{D}^0\) and \(B^0_{d,s} - \bar{B}^0_{d,s}\) systems.

For \(B_{d,s}\) mesons, in the SM the short distance contributions are expected to be the dominant ones. The prediction for \(\Delta \Gamma\) for \(B^0_s - \bar{B}^0_s\) is \[8\]  
\[
\Delta \Gamma_s = (0.106 \pm 0.032) \text{ ps}^{-1}.
\]

(1.3)
which gives \( y_{SM} = 0.078 \pm 0.025 \). The DØ experiment has measured this width difference \( \Delta \Gamma_d \) (see also \( \ref{eq:1} \)). Allowing the non-zero CP violation in mixing they obtained, \( \Delta \Gamma_s = (0.17 \pm 0.09_{\text{stat}} \pm 0.03_{\text{syst}}) \text{ ps}^{-1} \) \( y = 0.125 \pm 0.066_{\text{stat}} \pm 0.022_{\text{syst}} \), and in the CP conserving limit, \( \Delta \Gamma_s = (0.12 \pm 0.08_{\text{stat}}^{+0.03}_{-0.04_{\text{syst}}}) \text{ ps}^{-1} \) \( y = 0.088 \pm 0.059_{\text{stat}}^{+0.022}_{-0.030_{\text{syst}}} \). Within error bars, SM agrees with data. However, it is interesting to see whether the width difference can be much larger than the SM.

For the \( B_d \rightarrow \bar{B}_d \) system the width difference in SM is known to be small \( \Delta \Gamma_d = (26.7 \pm 0.08) \times 10^{-4} \text{ ps}^{-1} \), corresponding to \( y_{SM} = (2.058 \pm 0.006) \times 10^{-3} \). There is no experimental data for the width difference yet. It is interesting to see whether the width difference can be larger than the SM.

For the \( D^0 - \bar{D}^0 \) mixing in the SM as was shown in \( \ref{eq:2} \) \( x \) and \( y \) are generated only at the second order in SU(3) breaking, \( x, y \sim \sin^2 \theta_C \times [\text{SU(3) breaking}]^2 \). Most of the studies give \( x, y < 10^{-3} \), although large values are not excluded \( \ref{eq:3} \).

Recently, the parameter \( y \) for \( D^0 - \bar{D}^0 \) mixing has been measured. BaBar, assuming no CP violation in mixing, has analyzed the doubly Cabibbo suppressed (DCS) \( D^0 \rightarrow K^+ \pi^- \) mode \( \ref{eq:4} \), while Belle has studied singly Cabibbo suppressed (SCS) \( D^0 \rightarrow K^+ K^-, \pi^+ \pi^- \) decays \( \ref{eq:5} \). From these results the authors in \( \ref{eq:6} \) have fitted the mixing parameters and get the following result for \( y \) with 68\% and 95\% probability correspondingly

\[
y = (6.1 \pm 1.9) \times 10^{-3}, \quad y \in [0.0023, 0.0102].
\] (1.4)

In this work, we find that R-parity violating contribution to the parameter \( y \) is small for \( D^0 - \bar{D}^0 \) system, less than a few times \( 10^{-4} \). The RPV contribution for \( B_d \rightarrow \bar{B}_d \) system can be larger than the SM prediction. For \( B_s \rightarrow \bar{B}_s \) system, the contribution to \( y \) can be as large as the SM contribution.

In the following sections, we provide the detailed calculations.

### 2. General expression for \( \Gamma_{12} \)

Before going into specific RPV model calculations, we summarize some general results for short distance NP contribution to \( \Gamma_{12} \) from four quark operators generated by SM and NP. The calculation is straightforward. Starting from tree level four quark interactions, one needs to obtain the absorptive part for figure 1. Let us take \( D^0 - \bar{D}^0 \) mixing for illustration. For the cases considered here, we can write the \( \Delta C = -1 \) interaction Lagrangian as

\[
\mathcal{L}^{\Delta C = -1} = \sum_{q,q'} \left\{ \mathbf{D}_{qq'} [C_1(\mu)Q_1 + C_2(\mu)Q_2] + \mathbf{D}'_{qq'} [C'_1(\mu)Q_3 + C'_2(\mu)Q_4] \right\}, \tag{2.1}
\]

\[
Q_1 = \bar{u}_i \Gamma_1 q'_j \bar{q}_j \Gamma_2 c_i, \quad Q_2 = \bar{u}_i \Gamma_1 q'_i \bar{q}_j \Gamma_2 c_j, \quad Q_3 = \bar{u}_i \Gamma_3 q'_j \bar{q}_j \Gamma_4 c_i, \quad Q_4 = \bar{u}_i \Gamma_3 q'_i \bar{q}_j \Gamma_4 c_j.
\]

In the above we have omitted possible Lorentz indices for \( \Gamma_i \) which are contracted. The specific form of \( \Gamma_i \) depends on the nature of interaction generating the four quark operators. The notations here are that \( \Gamma_{1,2(3,4)} \) and \( \Gamma_{3,4(1,2)} \) should appear on the left and right four quark vertices in figure 1, respectively.
Evaluating the diagram in figure 1, one obtains the following general expression for $\Gamma_{12}$:

$$\Gamma_{12} = -\frac{1}{2M_D} \sum_{q,q'} D_{qq'} D'_{q'q} (K_1 \delta_{ik} \delta_{jl} + K_2 \delta_{ii} \delta_{jk}) \sum_{\alpha=1}^{5} I_{\alpha}(x,x') \langle D^0 | O_{ij}^{\alpha \ell \ell} | D^0 \rangle,$$

where

$$K_1 = (C_1 C'_1 N_c + (C_1 C'_2 + C'_1 C_2)) \quad K_2 = C_2 C'_2. \quad (2.3)$$

The operators are defined as

$$O_1^{ij \ell \ell} = \bar{u}_k \Gamma_3 \gamma_\mu \Gamma_2 c_j \bar{u}_\ell \Gamma_1 \gamma_\nu \Gamma_4 c_i,$$
$$O_2^{ij \ell \ell} = \bar{u}_k \Gamma_3 \gamma_\mu \Gamma_2 c_j \bar{u}_\ell \Gamma_1 \gamma_\nu \Gamma_4 c_i,$$
$$O_3^{ij \ell \ell} = \bar{u}_k \Gamma_3 \gamma_\mu \Gamma_2 c_j \bar{u}_\ell \Gamma_1 \gamma_\nu \Gamma_4 c_i,$$
$$O_4^{ij \ell \ell} = \bar{u}_k \Gamma_3 \gamma_\mu \Gamma_2 c_j \bar{u}_\ell \Gamma_1 \gamma_\nu \Gamma_4 c_i,$$

and the coefficients $I_{\alpha}(x,x')$ are given by

$$I_1(x,x') = -\frac{k^* m_c}{48 \pi} \left[ 1 - 2(x + x') + (x - x')^2 \right],$$
$$I_2(x,x') = -\frac{k^*}{24 \pi m_c} \left[ 1 + (x + x') - 2(x - x')^2 \right],$$
$$I_3(x,x') = \frac{k^*}{8 \pi} \sqrt{x'} (1 + x' - x),$$
$$I_4(x,x') = -\frac{k^*}{8 \pi} \sqrt{x'} (1 - x' + x),$$
$$I_5(x,x') = \frac{k^* m_c}{4 \pi} \sqrt{x}, \quad (2.4)$$

where $k^* \equiv (m_c/2)[1 - 2(x + x') + (x - x')^2]^{1/2}$ with $x \equiv m_q^2/m_c^2$ and $x' \equiv m_{q'}^2/m_c^2$. Replacing $D'_{qq'}$ with the SM couplings and $\Gamma_3 = \gamma_\mu (1 - \gamma_5)/2$ and $\Gamma_4 = \gamma_\mu (1 - \gamma_5)/2$, we obtain the formula presented in [12, 13] for SM and NP interference contribution. Note that when considering the contributions from the same operator, one should take $D'_{qq'} = D_{qq'}$, $\Gamma_{1,2} = \Gamma_{3,4}$ and the expression for $\Gamma_{12}$ should be divided by 2.

Using the above formula, one can easily work out the expressions contributing from the SM (taking SM operators for $Q_{1,2,3,4}$), the interference between the SM and NP (SM...
- NP) (taking $Q_{1,2}$ from SM (NP) and $Q_{3,4}$ from NP (SM)), and purely NP (NP - NP) (taking $Q_{1,2,3,4}$ from NP). New physics effects can show up in the later two cases. We will concentrate on these contributions.

We comment that the fermions in the loop are not necessary to be quarks. They can be leptons too. If one identifies $q$ and $q'$ to be leptons, the correct result can be obtained by setting $N_c = 1$ and $C_{2}^{0} = 0$ in eq. (2.3). One can easily generalize the above formula for $B_{d,s}^{0} - ar{B}_{d,s}^{0}$ mixing cases with appropriate replacement of quark fields and couplings.

3. RPV contributions to $\Gamma_{12}$ for $D^{0} - \bar{D}^{0}$ mixing

In this section we give expressions for contributions to $\Gamma_{12}$ with RPV interactions.

Contributions from $\lambda'$ interaction to $\Gamma_{12}$ are given by

$$
\Gamma_{12, (SM-RPV)} = \frac{\sqrt{2} G_F X_{V_{12}^{\prime}} V_{c s}^{\star} m_{c}^{2} (C_{1} + C_{2}) \langle Q \rangle ,}
$$

$$
\Gamma_{12, (RPV-RPV,l)} = \frac{\lambda_{V_{12}^{\prime}}^{2} X_{V_{12}^{\prime}}^{\star} (C_{1} + C_{2}) m_{c}^{2}}{192 \pi m_{D}^{2} m_{D}^{2} m_{D}^{2} m_{D}^{2} (1/2 \langle Q \rangle - \langle Q_{s} \rangle)} ,
$$

$$
\Gamma_{12, (RPV-RPV,d)} = \frac{\lambda_{V_{12}^{\prime}}^{2} X_{V_{12}^{\prime}}^{\star} (C_{1} + C_{2}) m_{c}^{2}}{192 \pi m_{D}^{2} m_{D}^{2} m_{D}^{2} m_{D}^{2} (1/2 \langle Q \rangle - \langle Q_{s} \rangle)} ,
$$

where

$$
\langle Q \rangle = \langle \bar{D}^{0} | \bar{u}_{\alpha} \bar{\gamma}_{\mu} P_{L} e_{\alpha} \bar{u}_{\beta} \gamma_{\mu} P_{L} c_{\beta} | D^{0} \rangle , \quad \langle Q_{s} \rangle = \langle \bar{D}^{0} | \bar{u}_{\alpha} P_{R} c_{\alpha} \bar{u}_{\beta} P_{R} c_{\beta} | D^{0} \rangle ,
$$

The first equation in eq. (3.1) is the leading order result in $x_{s}$. Depending on the internal lepton exchanges, in the expression for $\Gamma_{12, (RPV-RPV,l)}$ the indices $i,j$ take 1 and 2 indicating which charged leptons are in the loop. In principle, one can also have an electron and a tauon in the loop. However, the tauon mass is close to the D meson mass, the contribution is suppressed by phase space. We will neglect this contribution. In the expression for $\Gamma_{12, (RPV-RPV,d)}, j,j'$ take 1 and 2 indicating which of the down quarks are in the loop.

$\Gamma_{12, (SM-RPV)}$ comes from SM interaction with the second term, $\Gamma_{12, (RPV-RPV,l)}$ comes from the first term, and $\Gamma_{12, (RPV-RPV,d)}$ comes from the second term, in $\mathcal{L}_{\text{eff}}(\lambda')$, respectively. Note that the SM-RPV contribution is proportional to the internal quark masses and the dominant one comes from $s\bar{s}$ in the loop. This is due to the chiral structure of $\Gamma$, which allow only $O_{5}^{ijkl}$ to contribute and therefore proportional to the function $I_{5}(x,x')$. If $x$ or $x'$ takes the down quark mass $I_{5}(x,x')$ is negligibly small.

In obtaining the expression for $\Gamma_{12, (SM-RPV)}$, we have used the SM $\Delta C = -1$ Lagrangian,

$$
\mathcal{L}_{\text{SM}} = -\frac{4 G_{F}}{\sqrt{2}} V_{q_{c} q_{u}}^{\star} [C_{1}(m_{c}) Q_{1} + C_{2}(m_{c}) Q_{2}] ,
$$

with $\Gamma_{1}$ and $\Gamma_{2}$ in eq. (2.1) to be $\gamma_{\mu}(1 - \gamma_{5})/2$ and $\gamma_{\mu}(1 - \gamma_{5})/2$, respectively.
The contributions from $\lambda''$ interaction come from the first term in $\mathcal{L}_{\text{eff}}(\lambda'')$ and are given by

$$
\Gamma_{12(\text{SM-RPV})} = -\frac{\sqrt{2} G_F \lambda''_{1/2} \lambda''^*_{2/2} V_{u3} V_{c3}^*}{8 \pi m_D m^2_{d'}} x_s m^2_c (2 C_1 + C_2) \langle Q' \rangle - C_2 \langle Q' \rangle ,
$$

$$
\Gamma_{12(\text{RPV-RPV})} = \frac{\lambda''_{1} \lambda''^*_{2} \lambda''_{1} \lambda''^*_{2}}{192 \pi m^2_{d'} m^2_{d'} m_D} \left( \frac{3}{2} \langle Q \rangle \right) .
$$

where

$$
\langle Q' \rangle = \langle D^0 | \bar{u}_a \gamma^\mu P_L c_a \bar{\nu}_\beta \gamma^\mu P_R c_\beta | D^0 \rangle , \quad \langle \tilde{Q}' \rangle = \langle \tilde{D}^0 | \bar{u}_a \gamma^\mu P_L c_a \bar{\nu}_\beta \gamma^\mu P_R c_\beta | D^0 \rangle .
$$

In first equation of eq. (3.3), as in the first equation of eq. (3.1), we only kept the leading order in $x_s$. The SM-RPV contribution is dominated by $s \bar{s}$ pair in the loop for the same reason as that for the $\lambda'$ case for SM-RPV contribution explained earlier.

Here we should mention that recently in ref. [12] the authors have considered RPV with slepton and squark exchanges for SM-NP contributions. Our predictions in the first equations in eqs. (3.1) and (3.3), for the same measurable, do not agree with their eqs. (16) and (24), respectively.

4. RPV contributions to $\Gamma_{12}$ for $B_{d,s}^0 - \tilde{B}_{d,s}^0$ mixing

In this case all terms except the first term in $\mathcal{L}_{\text{eff}}(\lambda')$ contribute to $\Gamma_{12}$.

4.1 The $\lambda'$ contribution

The expressions for $\Gamma_{12}$ from various contributions are given by

$$
\Gamma_{12(\text{SM-RPV})} = \frac{\sqrt{2} G_F m^2_{\lambda'} \lambda''_{1/2} \lambda''^*_{2/2} V_{u3} V_{c3}^*}{48 \pi m_B m^2_{d'}} \left\{ (2 C_1 (m_b) - C_2 (m_b)) \langle Q' \rangle + (2 C_2 (m_b) - C_1 (m_b)) \langle Q' \rangle \right\} ,
$$

$$
\Gamma_{12(\text{RPV-RPV},u)} = \frac{m^2_b}{192 \pi m_B} \left\{ \frac{\lambda'_{jik} \lambda'_{j'k'}}{m^2_{d'}} \frac{\lambda'_{j'k}\lambda'_{j3}}{m^2_{d}} \langle Q \rangle + \langle Q_s \rangle \right. 

+ \frac{\lambda'_{jik} \lambda'_{j'k3}}{m^2_{d'}} \frac{\lambda'_{j'k}\lambda'_{j3}}{m^2_{d}} \langle Q \rangle + \langle Q_s \rangle 

- \left. \frac{2 \lambda'_{jik} \lambda'_{j'k3}}{2 m^2_{d'}} \frac{\lambda'_{j'k}\lambda'_{j3}}{m^2_{d}} \left( \langle Q' \rangle - \frac{1}{2} \langle Q' \rangle \right) \right\} ,
$$

$$
\Gamma_{12(\text{RPV-RPV},l)} = \frac{m^2_b}{192 \pi m_B} \left\{ \frac{\lambda'_{jik} \lambda'_{j'k3}}{m^2_{u'}} \frac{\lambda'_{j'k}\lambda'_{j3}}{m^2_{u}} \langle Q \rangle + \langle Q_s \rangle \right\} ,
$$

$$
\Gamma_{12(\text{RPV-RPV},u)} = \frac{m^2_b}{192 \pi m_B} \left\{ \frac{\lambda'_{jik} \lambda'_{j'k3}}{m^2_{u'}} \frac{\lambda'_{j'k}\lambda'_{j3}}{m^2_{u}} \left( \frac{1}{2} \langle Q \rangle - \langle Q_s \rangle \right) \right\} ,
$$
In this section, we carry out numerical analysis for RPV contributions to mixing parameter y for the SM, they have a simple relation
\[
\frac{\Gamma_{12}(RPV_{-RPV,d})}{\Gamma_{12}(RPV_{-RPV,u})} = \frac{m_\nu^2}{192\pi m_B} \frac{1}{m_{\nu_L}^2 m_{\nu_R}^2} \left\{ 18\lambda_{ij}^{\nu} \lambda_{i3}^{\nu} \lambda_{j3}^{\nu} \langle \bar{Q}' \rangle + (\lambda_{ij}^{\nu} \lambda_{i3}^{\nu} \lambda_{j3}^{\nu} + \lambda_{ij}^{\nu} \lambda_{i3}^{\nu} \lambda_{j3}^{\nu}) \left( \frac{1}{2} \langle Q \rangle - \langle Q_s \rangle \right) \right\},
\]
where \( j, j' \) take the values 1 and 2.

\[
\lambda_{cck} = V_{ck}^{\nu} V_{cb} \lambda_{12k}^{\nu}, \quad \lambda_{cuk} = V_{ck}^{\nu} V_{ub} \lambda_{11k}^{\nu},
\]
\[
\lambda_{uck} = V_{uk}^{\nu} V_{cb} \lambda_{12k}^{\nu}, \quad \lambda_{auk} = V_{uk}^{\nu} V_{ub} \lambda_{11k}^{\nu}.
\]

and

\[
\langle Q' \rangle = \langle B_0^q | q^k_\alpha \gamma_{\mu} P_L b_\alpha q^k_\beta \gamma_{\mu} P_R b_\beta | B_0^q \rangle,
\]
\[
\langle \bar{Q}' \rangle = \langle B_0^q | \bar{q}^k_\alpha \gamma_{\mu} P_L b_\alpha \bar{q}^k_\beta \gamma_{\mu} P_R b_\beta | B_0^q \rangle.
\]

Here for \( B_d \) and \( B_s \) systems, \( k \) takes 1 and 2, respectively.

The five different contributions to \( \Gamma_{12} \) listed above come from the second, first and sixth, fourth and fifth, second and third terms in \( \mathcal{L}_{\text{eff}}(\lambda') \), respectively.

4.2 The \( \lambda'' \) Contribution

In this case we have

\[
\Gamma_{12}(SM-RPV) = \frac{\sqrt{2}G_F x_e \sqrt{1 - 4x_c m_\nu^2 V_{ck}^{\nu} V_{cb} \lambda_{12k}^{\nu} \lambda_{123}^{\nu}}}{8\pi m_B m_\nu^2} \times \left\{ (2C_1(m_b) + C_2(m_b))\langle Q' \rangle - C_2(m_b)\langle \bar{Q}' \rangle \right\},
\]
\[
\Gamma_{12}(RPV-RPV_{,u}) = \frac{m_\nu^2}{192\pi m_B} \lambda_{ij}^{\nu} \lambda_{i3}^{\nu} \lambda_{i3}^{\nu} \lambda_{j3}^{\nu} \left( \frac{3}{2} \langle Q \rangle \right),
\]
\[
\Gamma_{12}(RPV-RPV_{,d}) = \frac{m_\nu^2}{48\pi m_B} \lambda_{ij}^{\nu} \lambda_{ij}^{\nu} \lambda_{i3}^{\nu} \lambda_{i3}^{\nu} \lambda_{i3}^{\nu} \left( \frac{3}{2} \langle Q \rangle \right),
\]

where \( i, i' \) take the values 1 and 2. The first two terms are due to the first term in \( \mathcal{L}_{\text{eff}}(\lambda'') \), and the last term is due to the second term in \( \mathcal{L}_{\text{eff}}(\lambda'') \).

The SM-RPV interference is dominated by \( c\bar{c} \) exchange in the loop for the same reasons as that for the \( s\bar{s} \) dominance for SM-RPV \( D^0 \rightarrow \bar{D}^0 \) mixing.

5. Numerical analysis

In this section, we carry out numerical analysis for RPV contributions to mixing parameter y for \( D^0 \rightarrow \bar{D}^0 \) and \( B_0^0 \rightarrow B_0^0 \) systems.

In general RPV contribution to \( \Gamma_{12} \) has CP violating phases associated with the new couplings. The relation of \( \Gamma_{12} \) and \( y \) is not trivial. If CP violating effects can be neglected which is the case for the SM, they have a simple relation

\[
y = \frac{\Gamma_{12}}{\Gamma}.
\]
\( \mu = 1.3 \text{ GeV} \quad \mu = 4.8 \text{ GeV} \quad \text{Masses (GeV)} \quad \text{Decay cons. (GeV)} \quad \text{Widths (GeV)} \)

<table>
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<th>( x_s )</th>
<th>( x_c )</th>
<th>( M_D )</th>
<th>( f_D )</th>
<th>( \Gamma_D )</th>
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<td>( f_{B_d} )</td>
<td>( \Gamma_{B_d} )</td>
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<td>-0.272</td>
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<tr>
<td>( C_2 )</td>
<td>( C_2 )</td>
<td>( M_{B_s} )</td>
<td>( f_{B_s} )</td>
<td>( \Gamma_{B_s} )</td>
</tr>
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<td>1.208</td>
<td>1.120</td>
<td>5.368</td>
<td>0.260</td>
<td>( 4.46 \times 10^{-13} )</td>
</tr>
</tbody>
</table>

**Table 1:** The central values of input parameters and coefficients.

In our numerical analysis, we assume CP conservation for easy comparison with data and other constraints obtained in the literature.

To compare with data, one needs to evaluate various hadronic matrix elements in the expressions for \( \Gamma_{12} \). We write them in the following form

\[
\langle Q \rangle = \frac{2}{3} f_R^2 m_R^2 B_Q, \quad \langle Q_s \rangle = -\frac{5}{12} f_R^2 m_R^2 B_{Q_s},
\]

\[
\langle Q' \rangle = -\frac{5}{6} f_R^2 m_R^2 B_{Q'}, \quad \langle \tilde{Q}' \rangle = -\frac{7}{6} f_R^2 m_R^2 B_{\tilde{Q}'},
\]

(5.2)

where \( B_Q \) factors are the so called bag parameters \([4]\). This way of parameterizing the matrix elements was inspired by vacuum saturation approximation. In the vacuum saturation approximation, they are all equal to one, which we will use in our estimate.

In the table 1 we list the parameters and coefficients appearing in the equations above. The input CKM elements are \([16]\)

\[
\lambda \equiv |V_{us}| = 0.2248, \quad A\lambda^2 \equiv |V_{cb}| = 41.5 \times 10^{-3}.
\]

(5.3)

The charm quark mass also comes into the calculations. In our numerical analysis we identify \( m_c \) and \( m_b \) with pole masses. Numerically we use \([14]\) \( m_{c,\text{pole}}/m_{b,\text{pole}} = \sqrt{x_c} = 0.29 \), which is based on the mass difference \( m_{b,\text{pole}} - m_{c,\text{pole}} = 3.4 \text{ GeV} \) and \( m_{b,\text{pole}} = 4.8 \text{ GeV} \).

To give some understanding of RPV contributions, in the following analysis we take the central values for the input parameters.

5.1 \( D^0 - \bar{D}^0 \) mixing

Taking all \( \lambda' \) and \( \lambda'' \) to be real, and inserting known values for the parameters involved, for \( \lambda' \) contributions, we have

\[
y_{(SM-RPV)} = 0.0037 \times \frac{\lambda'_{22} \lambda'^*_{12} (100 \text{ GeV})^2}{m^2_{\tilde{e}_L}},
\]

\[
y_{(RPV-RPV,l)} = 0.3298 \times \frac{\lambda'_{2k} \lambda'^*_{1k} \lambda'_{2k} \lambda'^*_{1k} (100 \text{ GeV})^4}{m^2_{d_R} m^2_{d_R'}},
\]

\[
y_{(RPV-RPV,q)} = 0.9893 \times \frac{\lambda'_{2j} \lambda'^*_{1j} \lambda'_{2j} \lambda'^*_{1j} (100 \text{ GeV})^4}{m^2_{e_L} m^2_{e_L'}},
\]

(5.4)
The $\lambda''$ contributions are given by

$$y_{(SM-RPV)} = -0.0077 \times \lambda''_{1j2} \lambda''_{2j2} \frac{(100\text{GeV})^2}{m_{\tilde{g}}^2} ,$$

$$y_{(RPV-RPV)} = 1.3191 \times \lambda''_{1j1} \lambda''_{2j1} \lambda''_{1j2} \lambda''_{2j2} \frac{(100\text{GeV})^4}{m_{\tilde{g}}^2 m_{\tilde{e}_L}^2} .$$ (5.5)

There are constraints on the RPV parameters from various other processes [17, 19, 18]. Taking these constraints into account, we list in table 2 the corresponding values for the mixing parameter $y$.

<table>
<thead>
<tr>
<th>RPV parameters</th>
<th>Bounds [Processes]</th>
<th>Estimate</th>
</tr>
</thead>
<tbody>
<tr>
<td>$</td>
<td>\lambda''<em>{122}\lambda''</em>{112}</td>
<td>$</td>
</tr>
<tr>
<td>$</td>
<td>\lambda''<em>{12k}\lambda''</em>{11k}</td>
<td>$</td>
</tr>
<tr>
<td>$</td>
<td>\lambda''<em>{i2j}\lambda''</em>{j1j}</td>
<td>_{j=1,2}$</td>
</tr>
<tr>
<td>$</td>
<td>\lambda''<em>{132}\lambda''</em>{232}</td>
<td>$</td>
</tr>
<tr>
<td>$</td>
<td>\lambda''<em>{132}\lambda''</em>{232}</td>
<td>$</td>
</tr>
</tbody>
</table>

Table 2: The bounds on parameters from [17, 18, 19, 21] and corresponding values for $y$.

5.2 $B_d^0 - \bar{B}_d^0$ mixing

Here we present the numerical results for the $B_d^0 - \bar{B}_d^0$ mixing. For $\lambda'$ contributions, we have

$$y_{(SM-RPV)} = -224.5 \times \lambda'_{q1} \frac{(100\text{GeV})^2}{m_{\tilde{e}_L}^2} ,$$
We list various constraints on relevant RPV parameters and corresponding values for $y$.

<table>
<thead>
<tr>
<th>RPV parameters</th>
<th>Bounds [Processes]</th>
<th>Estimate</th>
</tr>
</thead>
<tbody>
<tr>
<td>$</td>
<td>\lambda''<em>{212}\lambda''</em>{313}</td>
<td>$</td>
</tr>
<tr>
<td>$</td>
<td>\lambda''<em>{212}\lambda''</em>{233}</td>
<td>$</td>
</tr>
<tr>
<td>$\lambda''<em>{ij}\lambda''</em>{ij3}$</td>
<td>$1.1 \times 10^{-3}$ $[B^0 \to X_q \nu \bar{\nu}]$</td>
<td>$y_{(RPV-RPV)_{ij3}} \approx 2.7 \times 10^{-4}$</td>
</tr>
<tr>
<td>$\lambda''<em>{ij}\lambda''</em>{ij3}$</td>
<td>$1.1 \times 10^{-3}$ $[B^0 \to X_q \nu \bar{\nu}]$</td>
<td>$y_{(RPV-RPV)_{ij3}} \approx 0.67 \times 10^{-4}$</td>
</tr>
<tr>
<td>$\lambda''<em>{ij}\lambda''</em>{ij3}$</td>
<td>$6.4 \times 10^{-7}$ $[B\bar{B}]$</td>
<td>$y_{(RPV-RPV)_{ij3}} \approx 1.1 \times 10^{-4}$</td>
</tr>
<tr>
<td>$\lambda''<em>{ij}\lambda''</em>{ij3}$</td>
<td>$1.6 \times 10^{-6}$ $[B\bar{B}]$</td>
<td>$y_{(RPV-RPV)_{ij3}} \approx 7.2 \times 10^{-4}$</td>
</tr>
<tr>
<td>$\lambda''<em>{i12}\lambda''</em>{323}$</td>
<td>$6 \times 10^{-5}$ $[B \to \phi \pi]$</td>
<td>$y_{(SM-RPV)_{i12}} \approx 2.8 \times 10^{-5}$</td>
</tr>
<tr>
<td>$\lambda''<em>{i12}\lambda''</em>{323}$</td>
<td>$6 \times 10^{-5}$ $[B \to \phi \pi]$</td>
<td>$y_{(RPV-RPV)_{i12}} \approx 0.8 \times 10^{-4}$</td>
</tr>
<tr>
<td>$\lambda''<em>{i12}\lambda''</em>{323}$</td>
<td>$6 \times 10^{-5}$ $[B \to \phi \pi]$</td>
<td>$y_{(RPV-RPV)_{i12}} \approx 3.2 \times 10^{-4}$</td>
</tr>
</tbody>
</table>

Table 3: The bounds on parameters from [17, 19, 22, 24, 27] and corresponding values for $y$. 

For $\lambda''$ contributions, we have

$$y_{(SM-RPV)} = 0.46 \times \lambda''_{212}\lambda''_{323} \frac{(100 \text{GeV})^2}{m^2_{d_R}},$$

$$y_{(RPV-RPV)_{u}} = 220.3 \times \lambda''_{212}\lambda''_{233}\lambda''_{213}\lambda''_{233} \frac{(100 \text{GeV})^4}{m^2_{d_R} m^2_{d_R}},$$

$$y_{(RPV-RPV)_{d}} = 881.3 \times \lambda''_{j12}\lambda''_{j32}\lambda''_{j13}\lambda''_{j32} \frac{(100 \text{GeV})^4}{m^2_{d_R} m^2_{d_R}}.$$
\( \lambda_{q'q} \). Using constraints from \[23\], we have \( |\lambda_{121}^{e} \lambda_{113}^{e}| \lesssim 1.2 \times 10^{-5} \) and \( |\lambda_{121}^{e} \lambda_{123}^{e}| \lesssim 5.0 \times 10^{-5} \), each gives \( y \approx 1 \times 10^{-4} \). This value is much less than the SM prediction.

For \( y_{(RPV−RPV,u)} \), using the constraints \( |\lambda_{ij1}^{e} \lambda_{j13}^{e}|, |\lambda_{ij3}^{e} \lambda_{j33}^{e}| \lesssim 1.1 \times 10^{-3} \), from \[19\] we find that the corresponding upper bounds: \( y_{(RPV−RPV,u)} \simeq 2.7 \times 10^{-4} \) and \( y_{(RPV−RPV,l)} \simeq 6.7 \times 10^{-5} \).

As for the contribution \( y_{(RPV−RPV,u)} \), there are four terms with \( j, j' \) take values 1 or 2. Taking the explicit constraints from ref. \[23\], \( |\lambda_{i11}^{e} \lambda_{113}^{e}| \lesssim 8.0 \times 10^{-4}, |\lambda_{111}^{e} \lambda_{223}^{e}| \lesssim 2.5 \times 10^{-3}, |\lambda_{121}^{e} \lambda_{113}^{e}| \lesssim 1.2 \times 10^{-5} \), and \( |\lambda_{121}^{e} \lambda_{223}^{e}| \lesssim 5.0 \times 10^{-5} \), we find that the dominant contribution is from the case \( j = j' = 1 \) which gives the upper bound \( y_{(RPV−RPV,u)} \simeq 1.1 \times 10^{-4} \).

In the same way for the contribution \( y_{(RPV−RPV,d)} \), taking the constraints from \[22, 23\], the dominant part is from \( \lambda_{ijj}^{e} \lambda_{ijj}^{e} \lambda_{ijj}^{e} \lambda_{ijj}^{e} \) term with \( j = j' = 1 \). We find the value for \( y_{(RPV−RPV,d)} \) can be as large as \( 7 \times 10^{-3} \). This is about three times larger than the SM contribution.

Contributions from \( \lambda'' \) are also constrained. ref. \[25\] considers the decay mode \( B^- \to \phi \pi^- \) and drives the upper bound \( \lambda''_{ij1} \lambda''_{ij2} < 6 \times 10^{-5} \). Taking the same bound for \( |\lambda''_{ij1} \lambda''_{ij2}|; j=2 \) and \( |\lambda''_{ij1} \lambda''_{ij2}|; j=1 \) under the assumption that there is no accidental cancellation, \( y_{(SM−RPV)} \) and \( y_{(RPV−RPV,u(d))} \) are constrained to be small as can be seen from table \[3\].

We conclude that if there is no accidental cancellations, for \( B^0_{s} − B^0_{d} \) mixing, R-parity contribution to \( y \) can be as large as \( 7 \times 10^{-3} \) which is about three times larger than the SM prediction. This value is still difficult to be measured experimentally. However, if there is accidental cancellation, \( y \) could be bigger. Careful measurement of \( y \) for \( B^0_{s} − B^0_{d} \) can provide valuable information about new physics beyond the SM.

### 5.3 \( B^0_{s} − B^0_{d} \) mixing

For \( \lambda' \) contributions, we have

\[
y_{(SM−RPV)} = -316.6 \times \lambda_{q'q} \frac{(100 \text{GeV})^2}{m_{\tilde{e}_L}},
\]

\[
y_{(RPV−RPV,u)} = 77.7 \times (100 \text{GeV})^4 \left\{ \frac{\lambda_{j3i}^{e} \lambda_{j2i}^{e} \lambda_{j3i}^{e} \lambda_{j2i}^{e}}{m_{d_{R}}^2 m_{d_{R}}^2} \right\} + \frac{\lambda_{j2i}^{e} \lambda_{j3i}^{e} \lambda_{j2i}^{e} \lambda_{j3i}^{e}}{m_{d_{R}}^2 m_{d_{R}}^2} + 2 \frac{\lambda_{j3i}^{e} \lambda_{j2i}^{e} \lambda_{j3i}^{e} \lambda_{j2i}^{e}}{m_{d_{R}}^2 m_{d_{R}}^2} \right\} + 2 \frac{\lambda_{j3i}^{e} \lambda_{j2i}^{e} \lambda_{j3i}^{e} \lambda_{j2i}^{e}}{m_{d_{R}}^2 m_{d_{R}}^2},
\]

\[
y_{(RPV−RPV,l)} = 77.7 \times \lambda_{j2i}^{e} \lambda_{j3i}^{e} \lambda_{j2i}^{e} \lambda_{j3i}^{e} \frac{(100 \text{GeV})^4}{m_{d_{R}}^2 m_{d_{R}}^2},
\]

\[
y_{(RPV−RPV,u)} = 233.1 \times \lambda_{j2i}^{e} \lambda_{j3i}^{e} \lambda_{j2i}^{e} \lambda_{j3i}^{e} \frac{(100 \text{GeV})^4}{m_{d_{R}}^2 m_{d_{R}}^2},
\]

\[
y_{(RPV−RPV,d)} = 233.1 \times \left( \lambda_{j3i}^{e} \lambda_{j2i}^{e} \lambda_{j3i}^{e} \lambda_{j2i}^{e} + \lambda_{j2i}^{e} \lambda_{j3i}^{e} \lambda_{j2i}^{e} \lambda_{j3i}^{e} - 28 \lambda_{j3i}^{e} \lambda_{j2i}^{e} \lambda_{j3i}^{e} \right) \frac{(100 \text{GeV})^4}{m_{d_{R}}^2 m_{d_{R}}^2}.
\]
If we simply add them together we will get constraints from ref. [23] we have

\[ \lambda \] 

values for the mixing parameter can get for each contribution

\[ \lambda \] 

Table 4: Upper limits on parameters from [17, 22, 24] and corresponding values for \( y \). The numbers in the brackets correspond to the case, when central values for \( y_{SM} \) and \( y_{exp} \) are used to put the constraints. For each number see the text for the explanation.

<table>
<thead>
<tr>
<th>RPV parameters</th>
<th>Bounds [Processes]</th>
<th>Estimate</th>
<th>Our bounds on RPV</th>
</tr>
</thead>
<tbody>
<tr>
<td>(</td>
<td>\lambda'<em>{223}\lambda'</em>{221}</td>
<td>)</td>
<td>( 8.2 \times 10^{-3} [B, B_s] )</td>
</tr>
<tr>
<td>(</td>
<td>\lambda'<em>{22}\lambda'</em>{12}</td>
<td>)</td>
<td>( 1.5 \times 10^{-3} [B \rightarrow X_{\nu}] )</td>
</tr>
<tr>
<td>(</td>
<td>\lambda'<em>{12}\lambda'</em>{23}</td>
<td>)</td>
<td>( 1.5 \times 10^{-3} [B \rightarrow X_{\nu}] )</td>
</tr>
<tr>
<td>(</td>
<td>\lambda'<em>{12}\lambda'</em>{23}</td>
<td>), ( y \neq 3 )</td>
<td>( 5.16 \times 10^{-2} [B, B_s] )</td>
</tr>
<tr>
<td>(</td>
<td>\lambda'<em>{12}\lambda'</em>{23}</td>
<td>), ( y \neq 3 )</td>
<td>( 1.01 \times 10^{-2} [B \rightarrow K\pi] )</td>
</tr>
<tr>
<td>(</td>
<td>\lambda'<em>{12}\lambda'</em>{23}</td>
<td>), ( y \neq 3 )</td>
<td>( 1.2 \times 10^{-3} [B^+ \rightarrow \pi^+ K^0] )</td>
</tr>
</tbody>
</table>

For \( \lambda'' \) contributions, we have

\[ y_{(SM-RPV)} = -2.9 \times \lambda''_{221} \lambda''_{231} \left( \frac{100 GeV}{m_{\lambda''}} \right)^2, \]

\[ y_{(RPV-RPV,\nu)} = 310.8 \times \lambda''_{112} \lambda''_{13} \lambda''_{13} \lambda''_{123} \left( \frac{100 GeV}{m_{\lambda''}} \right)^4, \]

\[ (5.10) \]

\[ y_{(RPV-RPV,\nu)} = 1243 \times \lambda''_{221} \lambda''_{231} \lambda''_{231} \left( \frac{100 GeV}{m_{\lambda''}} \right)^4. \]

We list the constraints on the RPV parameters from [17, 13] and the corresponding values for the mixing parameter \( y \) in table 4.

There are several terms contributing to \( y \) from \( \lambda' \). For \( y_{(SM-RPV)} \) case we again drop terms proportional to \( V_{ub} \), and have, \( \lambda_{\nu\nu}=A\lambda^2\lambda'_{222} \times \left| \lambda'_{123} + \lambda\lambda'_{13} \right| \). We are using constraints from ref. [23] we have \( |\lambda'_{123}'\lambda'_{122}| \lesssim 8.2 \times 10^{-3} \) and from ref. [24] \( \left| \lambda'_{133}'\lambda'_{222} \right| \lesssim 2.48 \times 10^{-3} \). The first term dominates and gives \( y_{(SM-RPV)} \simeq 0.1, \) which is of order of SM prediction \( y_{SM} \simeq 0.078 \) and may have measurable effect.

For \( y_{(RPV-RPV,\nu)} \), we have three contributions. For first and second contributions using the following conditions on RPV parameters \( |\lambda'_{322}'\lambda'_{221}|, |\lambda'_{122}'\lambda'_{23}| \lesssim 1.5 \times 10^{-3} [17] \), we get \( y_{(RPV-RPV,\nu)} \simeq 1.7 \times 10^{-4} \). For the last term, we obtain \( y_{(RPV-RPV,\nu)} \simeq 3.5 \times 10^{-4} \). If we simply add them together we will get \( y_{(RPV-RPV,\nu)} \simeq 7 \times 10^{-4} \).

For \( y_{(RPV-RPV,d)} \), the situation is the same as the second term of \( (RPV-RPV,\nu) \) case.

In the case for \( y_{(RPV-RPV,\nu)} \), if one uses the individual constraints from [14, 23, 24] \( \left| \lambda'_{112}'\lambda'_{13}, \lambda'_{122}'\lambda'_{123}, \lambda'_{122}'\lambda'_{13} \right| = (1.63 \times 10^{-3}, 8.2 \times 10^{-3}, 5.16 \times 10^{-2}, 2.48 \times 10^{-3}) \) can get for each contribution \( y_{(RPV-RPV,\nu)} \simeq (6.2 \times 10^{-4}, 1.6 \times 10^{-2}, 6.0 \times 10^{-2}) \). If we keep the dominant interference term we will get \( y \simeq 0.06 \).

For \( y_{(RPV-RPV,d)} \) case we have three contributions. The contribution of the first term into \( y \) is small, about \( 1.6 \times 10^{-3} \). The dominant contributions here are coming from squares of \( \left| \lambda'_{322}'\lambda'_{221} \right| \simeq 1.29 \times 10^{-3} [B^- \rightarrow K^- \pi_0] \) and \( \left| \lambda'_{322}'\lambda'_{221} \right| \simeq 2.3 \times 10^{-3} [B^0 \rightarrow M \bar{M}] \) [13]. The second term is just the same as in \( y_{(RPV-RPV,\nu)} \) case, considered above. So here we have \( y_{(RPV-RPV,d)} \simeq 0.06 \). The last term is enhanced with the large coefficient.
Here the dominant contributions are coming from $|\lambda_{22}^{y} \lambda_{23}^{y} \lambda_{32}^{y} \lambda_{12}^{y}| \simeq 18.9 \times 10^{-6}$ and $|\lambda_{12}^{y} \lambda_{13}^{y} \lambda_{32}^{y} \lambda_{12}^{y}| \simeq 20.6 \times 10^{-6}$ \([23, 24, 10, 22]\). If we simply add them together, their contribution will be $y_{(RPV-RPV,d)} \simeq 0.26$. So, here we can conclude, that in $y_{(RPV-RPV,d)}$ case one also can expect large effects for $y$.

For $\lambda^{y}$ contribution to $y_{(SM-RPV)}$, we have $|\lambda_{221}^{y} \lambda_{231}^{y}| \lesssim 1.01 \times 10^{-2}$ from $B \rightarrow \bar{K} \pi$ \([24]\), which gives $y_{(SM-RPV)} \simeq 2.9 \times 10^{-2}$.

For the $y_{(RPV-RPV,u)}$ case no direct constraint on $|\lambda_{12}^{y} \lambda_{13}^{y}|$ exists. However, if one assumes that $|\lambda_{12}^{y} \lambda_{13}^{y}| \approx 1.2 \times 10^{-3}$ \([26]\), $y_{(RPV-RPV,v)} \simeq 4.5 \times 10^{-4}$.

For the last case of eq. (5.10) from \([26]\) we have $|\lambda_{12}^{y} \lambda_{13}^{y}| \lesssim 1.8 \times 10^{-3}$, which gives $y_{(RPV-RPV,d)} \approx 1.8 \times 10^{-3}$.

We note that present constraints on the RPV parameters still allow large $y_{(SM-RPV)}$, $y_{(RPV-RPV,u)}$ and $y_{(RPV-RPV,v)}$ from $\lambda^{y}$ interaction. One can turn the argument around to constrain the relevant RPV parameters by requiring that the new contributions do not exceed the allowed range for the difference of SM prediction and experimental values. We have carried out an analysis, taking the one sigma range values $y_{SM} = [0.054, 0.101]$ and $y_{Exp.} = [0.022, 0.151]$, and assumed constructive contributions between the SM and new contributions, to obtain the bounds for each individual terms. Similar analysis has been performed for $\lambda^{y}$ cases. The bounds are listed in table 4 in the last column. These bounds are new ones.

6. Conclusion

In this paper we have explored the influence of SUSY R-parity violation contributions for the lifetime difference $y$ on the $D^{0} - \bar{D}^{0}$ and $B_{u,s}^{0} - \bar{B}_{u,s}^{0}$ systems. We have obtained general expressions for new physics contributions to $y$ from effective four fermion operators including SM-NP interference and pure new physics contributions. We find that in general R-parity violating contribution to $D^{0} - \bar{D}^{0}$ mixing, and $B_{u,d}^{0} - \bar{B}_{d,u}^{0}$ to be small. There may be sizable contribution to $B_{u}^{0} - \bar{B}_{u}^{0}$ mixing. We also obtain some interesting bounds on R-parity violating parameters using known Standard Model predictions and experimental data.

Acknowledgments

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References


BELLE collaboration, M. Staric, Measurement of $D^0$-$\bar{D}^0$ mixing in $D^0 \to K^0_\pi^+\pi^-$ decays, talk given at the XLII Rencontres de Moriond, La Thuile, Italy, 10-17 March (2007) arXiv:0704.1009.


