Electroluminescence From the Ge Quantum Dot MOS Tunneling Diodes


Abstract—A Ge quantum dot (QD) light-emitting diode (LED) is demonstrated using a MOS tunneling structure for the first time. The oxide film was grown by liquid phase deposition at 50 °C to reduce the thermal budget. The infrared emission of ∼1.5 µm was observed from Ge QD MOS LEDs, similar to the p-type–intrinsic–n-type structure reported previously. At the negative gate bias, the electrons in the Al gate electrode tunnel to the Ge QD through the ultrathin oxide and recombine radiatively with holes to emit the ∼1.5-µm infrared. The electrons also recombine with holes in the Si cap, and the band edge emission from Si is also observed.

Index Terms—Ge quantum dot (QD) light-emitting diodes (LEDs), MOS tunneling diode.

I. INTRODUCTION

IT IS a long-sought goal to integrate ultralarge-scale integrated (ULSI) circuits with electrooptics to possibly overcome the speed limitation of electrical interconnects and to add extra functionalities on Si chip. Light-emitting diodes (LEDs) and detectors are essential devices to achieve this goal. The Si MOS LED has been demonstrated [1] to emit the 1.1-µm infrared at the Si band edge. The addition of Ge quantum dot (QD) into Si can tune the optical characteristic to longer wavelength, and the Ge QD MOS detector [2], [3] was reported to detect 1.3- and 1.5-µm infrared signals. The ∼1.5-µm infrared light emission was also observed in a previous study of electroluminescence (EL) from the Ge QD p-type–intrinsic–n-type (p-i-n) diodes [4]. However, the p-i-n structure is not fully compatible with the ULSI process. The building device in ULSI is the MOS structure. In this letter, we report that the Ge QD MOS tunneling diode can emit both the ∼1.5-µm infrared from the Ge QD and the ∼1.1-µm infrared from Si. The MOS tunneling diode reported here has the same structure used in the ULSI circuits, whereas the oxide is thin enough to have significant tunneling current from the metal to the substrate. The fabrication for MOS LED does not need the n-type doping in the ultrahigh-vacuum chemical vapor deposition (UHV/CVD).

II. DEVICE FABRICATION

The SiGe QDs are prepared by UHV/CVD on p-type Si (001) substrates with resistivities in the range of 15–25 Ω·cm. After an Si buffer layer of 100 nm was grown, 20 periods of Ge QD structure with base and height of ∼120 and ∼7 nm, respectively, were grown at a temperature of 600 °C under the Stranski–Krastanov (SK) growth mode [5] as shown in Fig. 1(a). The Ge layers are separated by 60-nm Si spacer layers. The Ge dot fabricated by the SK growth mode can lead to an increase of the luminescence efficiency because of the localization of excitons in the dots [6] and the increase of hole concentration in the dots. The area density of Ge islands is about 4 × 10^9 cm^{-2}. A 3-nm (nominal thickness) Si cap was deposited above the top layer of the self-assembled Ge layer as a cap layer for the subsequent deposition of liquid phase deposition (LPD) oxide. The undoped Si/Ge layers has a background hole concentration of ∼1 × 10^{16} cm^{-3}. Al is
used for the gate electrode and the back contact. Fig. 1(b) shows the device structure. To avoid material degradation such as strain relaxation and Ge out diffusion, the low-temperature oxide (LPD oxide) is used to reduce the thermal budget of the Si/Ge device process [3]. The tunneling current of electrons at the negative bias can cause electron–hole recombination at the oxide/Si interface and in the Ge QDs to emit at the negative bias can cause electron–hole recombination at the oxide/Si interface and in the Ge QDs to emit

**III. RESULTS AND DISCUSSION**

Fig. 2 shows Raman spectra excited by the 514-nm Ar laser for the Ge QD, bulk Ge, and bulk Si. In addition to the strong Si substrate signal at 520 cm\(^{-1}\), Ge-G0, Si-Ge, and local Si-Si vibrational peaks were also observed at \(\sim 302, \sim 417,\) and \(\sim 436\) cm\(^{-1}\), respectively. The appearance of the Si-Ge and the local Si-Si vibrational peaks implies the formation of Si-Ge alloy in the wetting layers [9]. On the other hand, the Ge-Ge peak is most possibly contributed by the islands [9]. As compared with the bulk Ge, the Raman shift of the Ge/Ge phonon frequency of the Ge QD with respect to the bulk Ge peak can be determined (\(\sim 2\) cm\(^{-1}\)). The possible factors responsible for this frequency shift are the alloying of the Ge QDs with Si spacers [10], the phonon confinement effect in the QDs, and the effects of compressive strain [11]. The first two factors reduce the phonon frequency, but the last mechanism increases the phonon frequency. Energy dispersion spectrometry shows the Ge dots transform into Si\(_{0.40}\)Ge\(_{0.54}\) dots due to Si/Ge interdiffusion at the growth temperature of 600 °C.

Fig. 3 shows the photoluminescence (PL) spectra at different temperatures from the Ge QD before the fabrication of the MOS LED device at a pump power of 60 W/cm\(^2\). There are two PL peaks located at \(\sim 1.1\) and \(\sim 0.85\) eV. These two peaks are most likely from the optical transitions related to the Si substrate and the Ge QD [12]. Fig. 4 shows the EL spectra at different temperatures from the Ge QD MOS LED samples with a device size of \(4 \times 10^{-4}\) cm\(^2\). The drive current is 100 mA at a gate voltage of \(-8\) V. Device temperatures of \(\sim 340, \sim 250, \sim 160,\) and \(\sim 110\) K are obtained from the fitting of the EL line shapes from the Si signal with the electron–hole–plasma recombination model [1], [7], whereas the temperatures of the cold finger in the EL measurement are 300, 200, 100, and 10 K, respectively. The temperature difference is due to the device heating. Because our Ge dot has a wide base (\(\sim 120\) nm) and a short height (\(\sim 7\) nm), the QD can be approximated by the quantum well in the electron–hole–plasma recombination model [13]. The band diagram of the Ge QD MOS LED at the accumulation condition is shown in the inset of Fig. 4. The traps in the LPD oxide and the small thickness (\(\sim 2\) nm) allow the significant tunneling electrons to inject from the Al electrode into the semiconductor. At a negative bias, the electrons tunnel from the Al gate to the substrate. Meanwhile, the negative gate bias also attracts holes in the p-type semiconductor and QDs, and the tunneling electrons can recombine with holes at the oxide/Si interface and in the Ge QD to emit \(\sim 1.1\) - and \(\sim 1.5\)-\(\mu\)m infrared, respectively [Fig. 1(b)]. The EL spectra of similar Ge QD p-i-n diode are also shown for reference. The EL spectra of similar Ge QD p-i-n diode are also shown for reference. The infrared emission of \(\sim 1.1\) - and \(\sim 1.5\)-\(\mu\)m observed from the Si and the Ge QD.
is a lower bound of the true quantum efficiency. The efficiency is too low to have product applications, but the EL of the MOS structure is a good characterization tool for material quality and can be used in the in-line inspection in the industry. Note that the MOS structure is much simpler than the p-i-n structure to measure the EL spectra.

IV. CONCLUSION

The ∼1.5-µm Ge QD MOS LED, which is fully compatible with the ULSI process, is reported for the first time. The origin of the emission is due to the radiative recombination between the electrons and the holes confined in the Ge QD. The electrons also recombined with the holes at the Si/oxide interface, and the band edge light emission from Si is also observed.

ACKNOWLEDGMENT

The authors would like to thank Prof. S. T. Chang and Prof. C.-T. Chia for helpful discussion.

REFERENCES


